



## **Overview of Methods**

- LASER PULSE SHAPING (for photocathodes)
  - $\rightarrow$  ATF @ BNL / UCLA (Zhou, Ben-Zvi, Babzien, Chang, Doyuran, Malone, Wang, Yakimenko)
  - $\rightarrow$  Sumitomo Heavy Industries Ltd. / Femtosecond Technology Research Association
- SOLENOIDS
  - $\rightarrow$  TTF @ DESY (Zhang)
  - $\rightarrow$  PITZ @ DESY / TU Darmstadt (Cee, Krassilnikov, Setzer, Weiland)
- NON-LINEAR ELECTROSTATIC FIELDS
  - $\rightarrow$  Eindhoven University of Technology / Pulsar Physics (van der Geer, de Loos, Botman, Luiten, van der Wiel)



### LEG

## Laser Pulse Shaping (1)

- Simulation code: PARMELA
- Laser beam non-uniformity:
  - $\implies$  Non-linear space charge forces
  - $\implies$  Emittance growth
- Investigate spatial and temporal shaping!





# Laser Pulse Shaping (2)

### Spatial shaping:

• Cylindrical symmetry: 30% - 40% emittance growth for peak-to-peak variations (40% - 70%); "hollow beam" leads to 100% increase!



• Non-cylindrical symmetry: Masks (90% – 50% transmission efficiency) show an increase of emittance between 30% - 100%







 $\implies \epsilon_n^{RMS} = 1.2\pi \text{ mm} \cdot \text{mrad for 1 nC current and 9 ps FWHM laser pulse}$ 





# $\operatorname{Solenoids}(1)$

- Simulation code: MAFIA
- Key Idea: Solenoid confines beam to forming Brillouin flow and transfers particles' angular momentum to a focusing force at its exit ⇒ Focusing Force = Space Charge Repulsion

 $\Longrightarrow$  Thereafter further acceleration to  $\gamma\gg 1$ 

- $\epsilon_x = \sqrt{\overline{x^2} \ \overline{p_x^2} \overline{xp_x^2}} \longrightarrow \text{linear forces don't alter } \epsilon_{RMS}$
- Busch's Theorem:  $rP_{\phi} + \frac{q}{2\pi}\psi = const$   $(\psi = \oint B \ d\sigma \text{ is the flux})$





• TTF setup:



- -B = 0 at the Neumann boundary; necessary for Busch's Theorem, otherwise beam would become axial-confined flow (each particle rotates around local magnetic flux line instead of global precession around beamline axis)
- Between a and c Brillouin flow (particles rotate around axis)
- Solenoid fringe very weak at  $c \to$  rotating momentum turns into focusing momentum



- Compensation solenoid: B = 0 at cathode  $\rightarrow$  electrons have no azimuthal momentum at gun exit



-  $z \ll 1$ : Strong emittance growth due to nonlinear radial space charge forces to about ≈ 5 mm·mrad

100

z/mm

2. iris

200

250

150

1. iris

50

#### $- z \gg 1$ : Emittance determined by RF effects and iris fringe fields

⇒ Emittance with and without solenoid field similar, but different slope of phase space ellipse (with solenoid: converges, without solenoid: beam exits gun in divergent state)









Radial component of the electric self-field in the median plane of a short and long bunch with 1 mm diameter and identical density profiles











## Non-linear Electrostatic Emittance Compensation (4)

- Solution:
  - Detrimental effects due to these field non-linearities must cancel each other
  - Radial third-order component of the electrostatic accelerating field should minimize transverse RMS emittance
- $\implies$  Special diode geometry:



Spherical aberration of diode lens:  $E_r = E_{r,1}(z) r + E_{r,3}(z) r^3$ 

## Non-linear Electrostatic Emittance Compensation (5)



- Uniform acceleration: only space charge
- Anode opening is at 2 mm  $\rightarrow$  largest  $E_{r,3}$
- For  $z \leq 1 \text{ mm}$  and  $z \geq 4 \text{ mm}$  acc. field is almost uniform and therefore  $E_{r,3}$  similar
- $\gamma \beta_r = p_r / mc$  vs. r is governed by  $E_{r,3}$  $\implies \gamma \beta_r = \gamma \beta_{r,1}(z) \ r + \gamma \beta_{r,3}(z) \ r^3$
- Special diode geometry:  $\gamma \beta_{r,3} = 0$  for  $z \ge 3.5$  mm
- $z \leq 1$  mm: roughly uniform field
- $1 \le z \le 1.5$  and  $2.5 \le z \le 3$ : compensation due to negative sign of  $E_{r,3}$
- Over-compensation at  $z \ge 1.5$  mm
- Strong emittance blow-up at z = 2.3 mm because of large positive  $E_{r,3}$

 $\implies$  Resulting  $\epsilon_n^{RMS}$  lower than with uniform acc. field

LEG





exit of the diode structure

LEG





## References

- Emittance Growth Due to the Laser Non-Uniformity in a Photoinjector, F. Zhou, I. Ben-Zvi, M. Babzien, X. Y. Chang, A. Doyuran, R. Malone, X. J. Wang, V. Yakimenko, Proceedings of EPAC 2002, Paris, France
- Experimental Studies of Photocathode RF Gun with Laser Pulse Shaping, J. Yang, F. Sakai, T. Yanagida, M. Yorozu, Y. Okada, T. Nakajyo, K. Takasago, A. Endo, Proceedings of EPAC 2002, Paris, France
- Intuitive Description of Emittance Compensation and its Application to Beam Transport Lines, M. Zhang, Proceedings of APAC 1998, Tsukuba, Japan
- Beam Dynamics Simulations for the PITZ RF-Gun, R. Cee, M. Krassilnikov, S. Setzer, T. Weiland, Proceedings of EPAC 2002, Paris, France
- Nonlinear Electrostatic Emittance Compensation in kA, fs Electron Bunches, S. B. van der Geer, M. J. de Loos, J. I. M. Botman, O. J. Luiten, M. J. van der Wiel, Physical Review E, Volume 65, 046501

The transparancies of this talk can be found at http://www.simonleemann.ch/work/leg/stateoftheart